

THE GEORGE Electro-Optic Plasmon Modulators: Breaking Photonic Limits WASHINGTON UNIVERSITY The George Washington University, School of Engineering and Applied Science (SEAS)

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Active Material

A combination of both modulation ability and light-matter-interaction response are critical for increased efficiency and device performance. Here we examine the free carrier and dispersive index tuning ability of two emerging active materials (Graphene and Indium Tin Oxide (ITO)) and the way we use these materials to modulate a light signal.

Refractive Index can be compared to the more recognized electrical imedance equation. The the formula for the complex refractive index (\tilde{n}) also has two parts:

$$\widetilde{n} = n + ki$$
 $Z = R + j\omega$

The real part of the refractive index (n) indicates the speed of the light's propogating phase, while the imaginary part of the refractive index quantifies the optical absorption (i.e., loss) of the media. Changing the optical loss is the key mechanism for these proposed electro-optic modulators. We are able to acheive this by altering the electrical carrier density of the device via an applied voltage bias. The refractive index of a medium can be related to its complex relative dialetric constant, $\tilde{\epsilon_r}$.

$$\epsilon_{r} = \epsilon_{1} + t_{0}$$
Graphene
$$\varepsilon_{eff}(\mu_{c}) = 1 - \frac{\sigma_{v}}{j\omega\varepsilon_{0}} = 1 - \frac{\sigma_{g}}{j\omega\varepsilon_{0}\Delta}$$

$$\Delta = \text{effective thickness of graphene}$$

$$\sigma = \text{electrical conductivity}$$

$$\varepsilon_{0} = \text{permitivity of free space}$$

$$\omega = \text{light angular frequency}$$
Indium Tin Oxide
$$\varepsilon(\omega) = \varepsilon_{\infty} - \frac{\omega_{p}^{2}}{(\omega + 1)^{2}}; \quad \omega_{p}^{2} = \frac{n_{c}e^{2}}{1}$$

 $\varepsilon_0 m^*$

 $\varepsilon_{\infty} =$ long-angular-momentum-limit permitivity ε_0 =permitivity of free space $\omega =$ angular momentum (rad/s) γ =electron scattering rate

 $\omega(\omega + i\gamma)'$



Device Structure

The structure of both of our hybrid plasmonic EOMs consists of an SOI waveguide and an a graphene/ITO - SiO2 - Au stack integrated on top. This specific configuration of materials forms a plasmonic HP mode and the Metal-Oxide-Semiconductor (MOS) capacitor, with the accumulation layer occuring at the active material's interface when a voltage bias is applied between the Gold and Silicon. Notice that this synergistic design allows for the triple use of the metal contact, namely (1) to form the HP mode, (2) to act as a heat sink, and (3) as an electrical electrode transporting the electrical data to the EOM gate, thus enabling for an ultra-compact design. Moreover, the overall length of the MOS stack is a mere 0.5um and 1um for the graphene - EOM and ITO - EOM, respectively. This compact device size is acheivable due to (1) the active materials ability to dramatically change its extinction coefficient (the imaginary part of the refractive index), and (2) the good overlap between the MOS mode and the active material. Notice that the peak of the electric field intensity across the MOS mode coincides per design directly with the area of the active material layer. Thus, when an electrical voltage is applied across our MOS capacitor, it forms an accumulation layer at the graphene/ITO-SiO2 interface, which in turn increases the active material's carrier density and, consequently, raises the extinction coefficient, k. This k-increase controls the optical absorption (i.e., the OFF-state) via Bear's law: $T(V) = Te^{-\alpha L}$



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Graphene EOM

Using the values from the complex index of refaction of graphene, the MOS field distribution changes between the modulator absorbing OFF and light-through ON state were calculated using a numerical finite element solver (Comsol). Various geometrical parameters were swept within Comsol with the goal of producing an optimized device design that exhibited a high extinction ratio and low power consumption.

Voltage (V)



Because EOMs play an integral role in the conversion between the electrical links, factors such as the scalability, modulation performance, and power a metric to benchmark EOMs in order to clearly demonstrate the stark contrast

References

[1] R. F. Oulton, V. J. Sorger et. al "Nano-photonic confinement and transport in a hybrid semiconductor-surface plasmon waveguide" Nature Phot. 2, 496-500 (2008) [2] Q. Xu, B. Schmidt, S. Pradhan, M. Lipson "Micrometre-scale silicon electro-optic modulator" Nature 435, 325-327 (2005) [3] A. Liu, R. Jones, L. Liao, D. Samara-Rubio et al. "A high-speed silicon optical modulator based on a metal-oxide-semiconductor capacitor" Nature 427, 615-619 (2004) [4] C. Huang, R. J. Lamond, S. K. Pickus, Z. R. Li, V. J. Sorger, "A Sub-λ-Size Modulator Beyond the Efficiency-Loss Limit" IEEE Photonic Journal 5, 4, 202-211 (2013) [5] V. J. Sorger, D. Kimura, R.-M. Ma and X. Zhang, "Ultra-compact silicon nanophotonic modulator with broadband response" Nanophotonics 1, 1, 17-22 (2012) [6] J. P. Donnelly, H. A. Haus, and N. Whitaker, "Analytical expressions for effective indices in a three-waveguide coupler" IEEE J.Quantum Electron. 23, 401 (1987) [7] F. Lou, D. Dai, and L. Wosinski "Ultracompact polarization beam splitter based on a dielectric-hybrid plasmonic-dielectric coupler" Opt. Lett. 37(16), 3372–3374 (2012) [8] Jingyee Chee, Shiyang Zhu, and G. Q. Lo, "CMOS compatible polarization splitter using hybrid plasmonic waveguide" **Opt. Express** 20, No. 23, 5 (2012)

$$f = 1 - \frac{T(L, \alpha_{ON})}{T_0} = 1 - e^{-\alpha_{ON}L} = 1 - e^{\Lambda}(-\frac{4\pi\kappa_{ON}}{\lambda}L)$$
$$\frac{V_b = V_{OFF}}{(V_b = V_{ON})} = \frac{T(L, \Delta\alpha)}{T_0} \rightarrow \eta_{mod} = \frac{ER}{V_{pp}}$$